

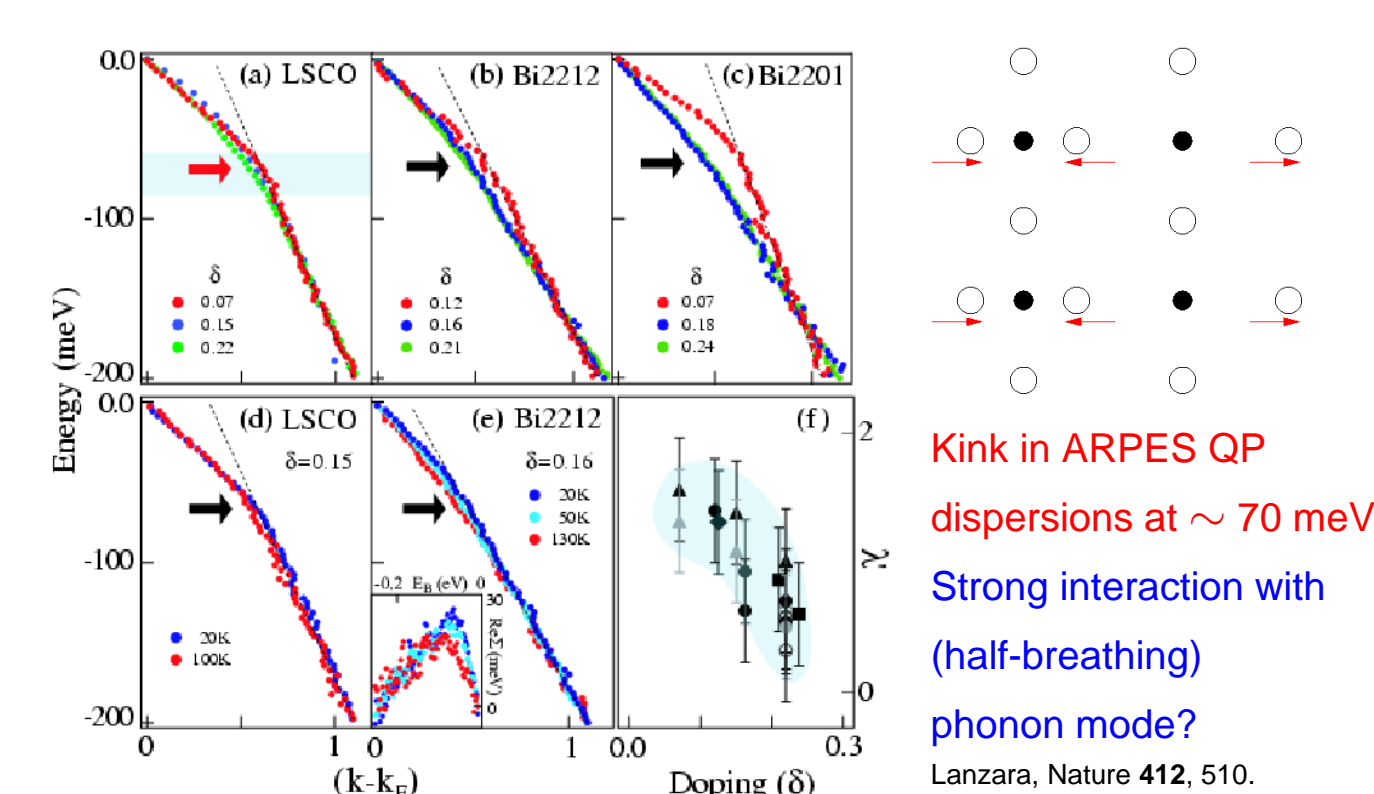
# Electron-phonon coupling to breathing phonons in cuprates

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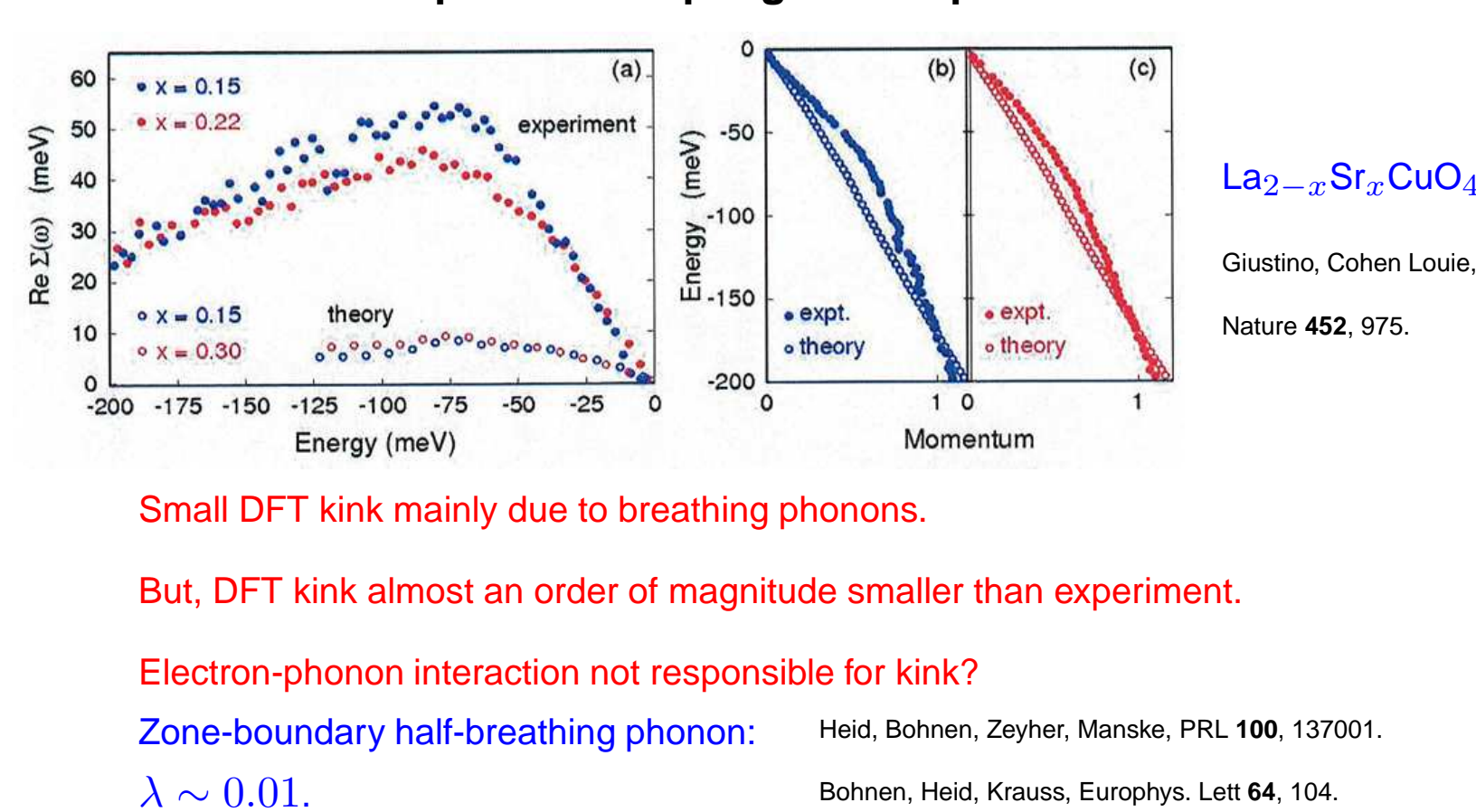
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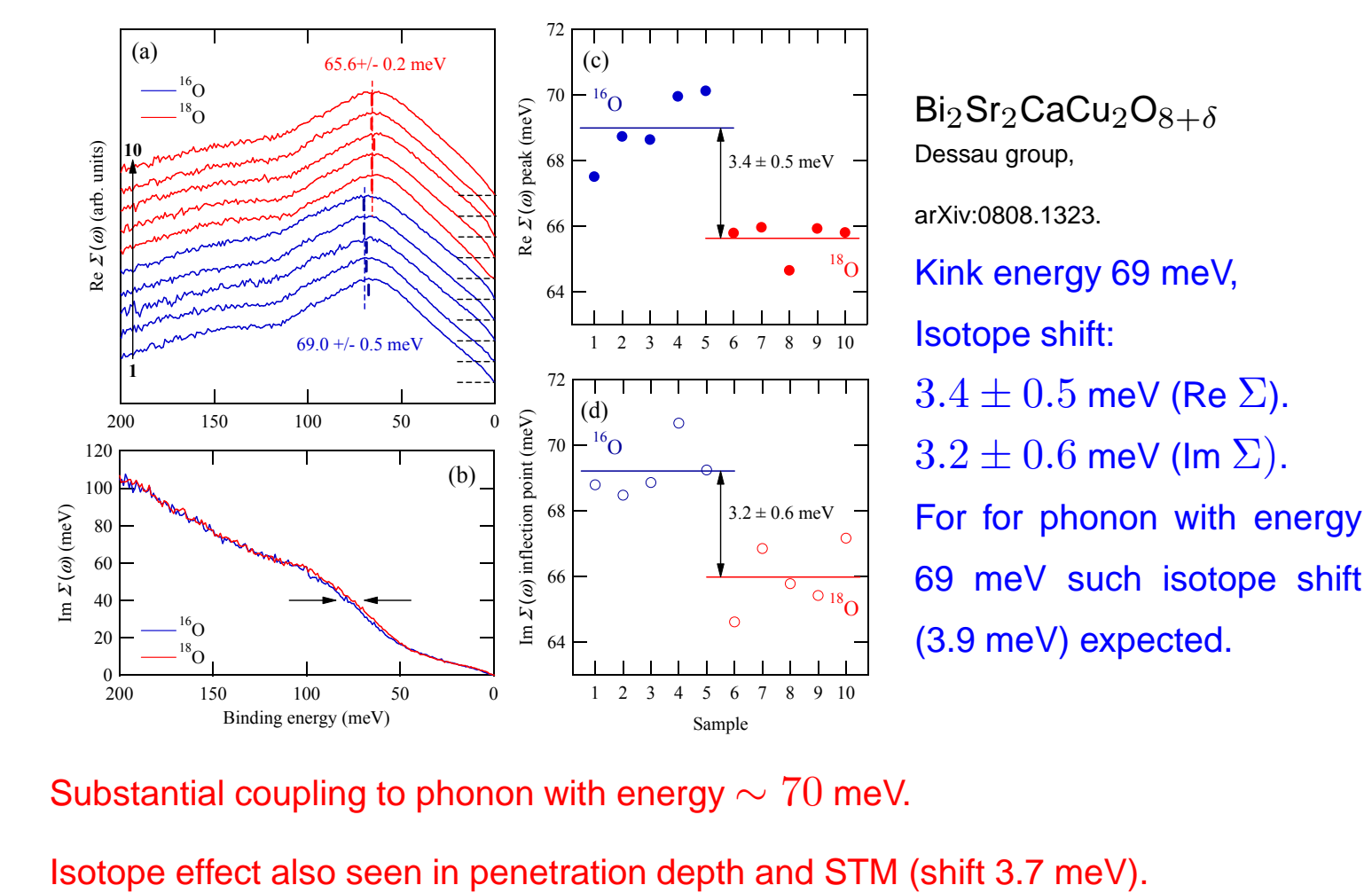
## Kink in electron dispersion



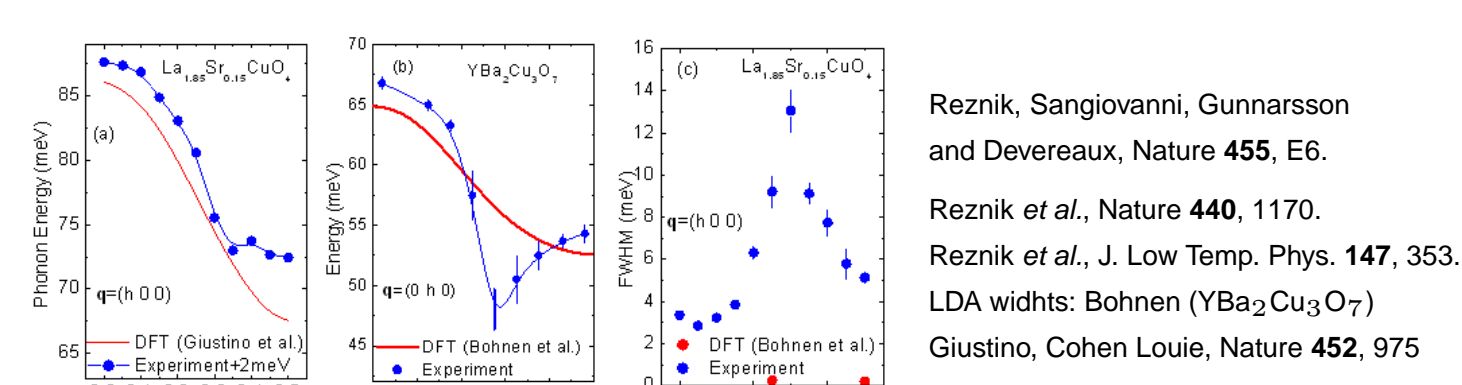
## DFT: Electron-phonon coupling and dispersion kinks



## Isotope shift of kink



## Half-breathing phonon. Neutron scattering



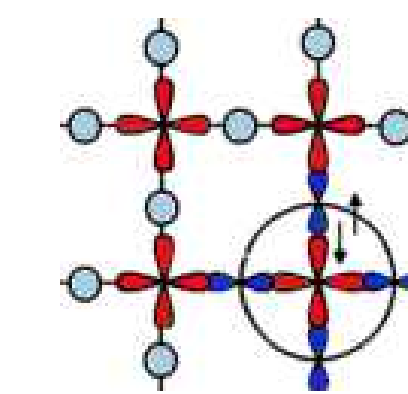
Very strong anomaly at  $q \approx (0.3, 0, 0)$  in both dispersion and width.  
Reduced with temperature  $\Rightarrow$  Not anharmonic effects or impurities.  
LDA usually describes phonons very well, but here strongly underestimates width.  
Anomalously strong coupling.

## Electron-phonon coupl. Undoped La2CuO4, Shell model

Find the electron-phonon coupling strength to a Zhang-Rice singlet.  
1. Use shell model to describe electrostatic coupling to phonons.  
Phonon eigenvectors  $\Rightarrow$  Potential on a singlet due to a phonon.  
Screening only by the "shells".  
LDA too effective screening.  
2. Add coupling due to modulation of  $t_{pd}$  and  $\epsilon_d - \epsilon_p$  (breathing phonons).  
Coupling to (half-)breathing modes (80 meV), O<sub>z</sub> modes (60-70 meV) and La (Cu) modes (20 meV).  
Total  $\lambda \sim 1.2$ . Breathing  $\lambda \sim 0.4$ .  
Sufficient to explain experimentally observed polaronic effects for undoped cuprates.  
Rösch, Gunnarsson, Zhou, Yoshida, Sasagawa, Fujimori, Hussain, Shen, Uchida, PRL 95, 227002.

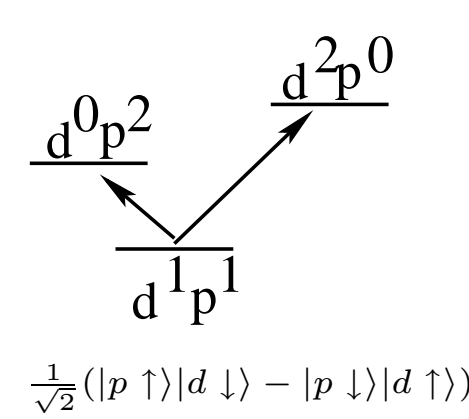
## Three-band (Emery) model

Important physics in CuO<sub>2</sub> plane.  
Include one Cu  $x^2-y^2-3d$  and two O  $2p$  orbital/cell.  
 $H = \sum_{i\sigma} \epsilon_i^d n_{i\sigma}^d + \sum_{j\sigma} \epsilon_j^p n_{j\sigma}^p + U \sum_i n_{i1}^d n_{i2}^d - \sum_{\langle i,j \rangle \sigma} t_{ij}^d (d_{i\sigma}^\dagger p_{j\sigma} + H.c.)$   
 $t$ - $J$  model  
Undoped system: One hole per Cu site. AF coupling.  
Doped system: Additional holes mainly on O sites.  
Form Zhang-Rice singlets from Cu and O holes.  
 $H_{eff} = J \sum_{\langle i,j \rangle} (\vec{S}_i \cdot \vec{S}_j - \frac{1}{4} n_i n_j) + \sum_{i,j,\sigma} t_{ij} \tilde{c}_{i\sigma}^\dagger \tilde{c}_{j\sigma}$



## Origin of strong electron-phonon coupling

Zhang-Rice singlet: Nominally  $d^1 p^1$ .  
Coherent hopping  $Cu x^2-y^2 \leftrightarrow O 2p \Rightarrow E^Z-R \approx -4t_{pd}^2 (\frac{2-\epsilon_p}{\epsilon_p} + \frac{2}{U-\epsilon_p})$ .  
Large singlet energy ( $\sim 5$  eV).  
Due to singlet character, coupling to two config.  
For rigid lattice uninteresting constant.  
Phonons modulate singlet energy via  $t_{pd}$ .  
Strong coupling because Zhang-Rice singlet energy large.  
Rösch and Gunnarsson, PRL 92, 146403.



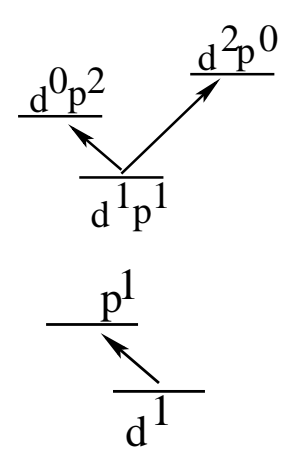
## HF Solution of three-band model

Consider three-band model with frozen phonons in Hartree-Fock (HF) mean-field approximation.  
Simulates LDA in paramagnetic doped system, but is able to give AF solution in undoped case  
Downfold to one-band model.  
Project out O-2p levels, get effective Cu-3d model.  
This allows direct comparison with  $t$ - $J$  model.

## Phonon softening in HF solution (linear response)

Softening given by  $Re \Pi_q(\omega) \approx g_q^2 \times \frac{1}{\omega} \int d\omega' \frac{Im \chi_q(\omega')}{\omega - \omega'}$   
 $\Pi_q(\omega) \approx g_q^2 \chi_q(\omega)$ : phonon self-energy  
 $\chi_q(\omega)$ : charge correlation function

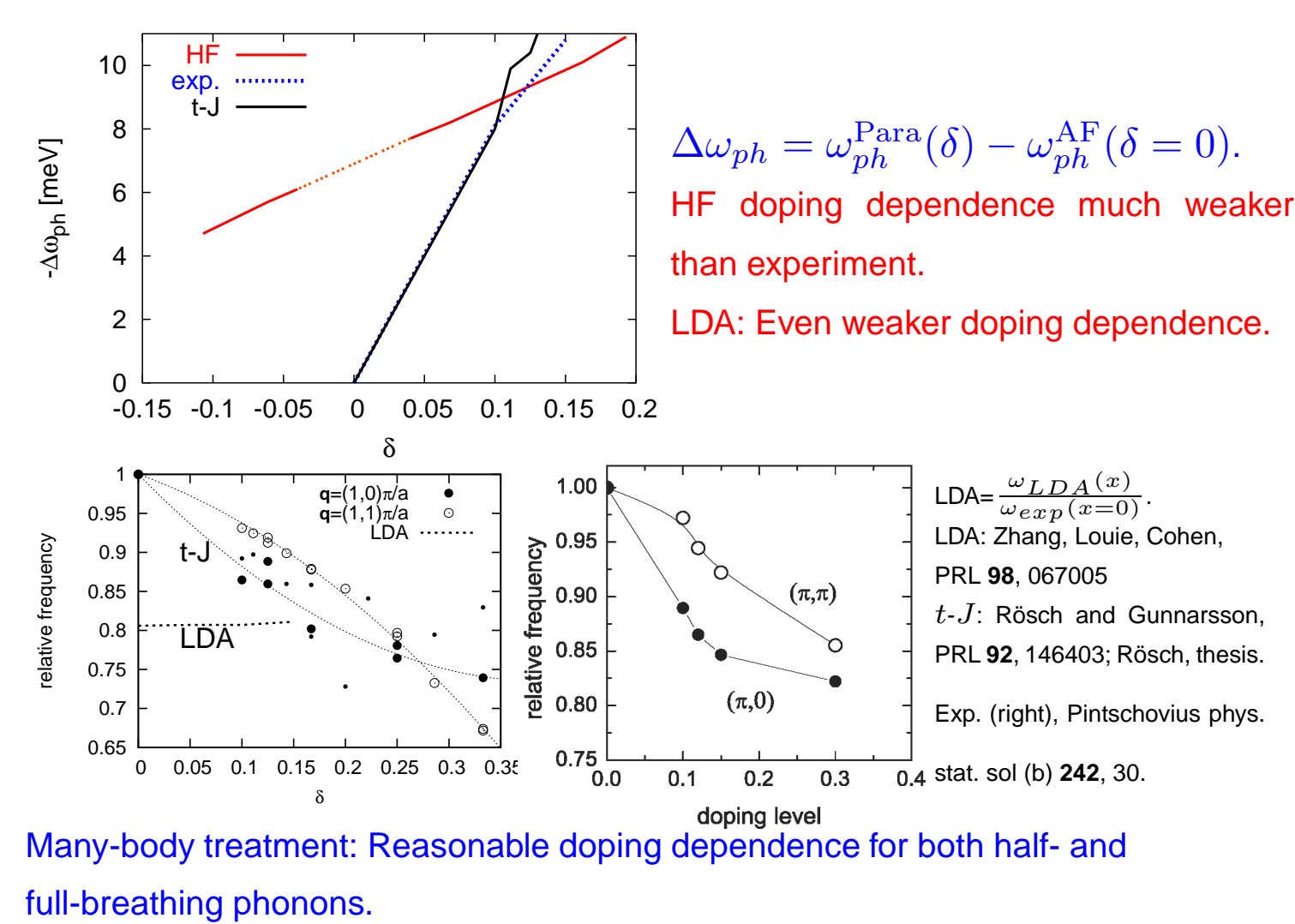
source	$t$ - $J$	HF	ratio
coupling $g^2$	$[\frac{2-\epsilon_p}{\epsilon_p} + \frac{2}{U-\epsilon_p}]^2$	$\frac{1}{\epsilon_p}$	$\approx 3$
sum rule for $Im \chi$	$\approx 2\pi\delta$	$\approx \pi$	$\approx 2\delta$
$< 1/\omega > Im \chi$	$\approx 1$	$\approx 0.5$	$\approx 1$
screening	$\approx 1$	$\approx 0.5$	$\approx 2$
product			$\approx 12\delta$



$\Rightarrow$  Softening for  $\approx 10 - 15\%$  doping is of right size but due to different reasons than in  $t$ - $J$  model

O. Rösch and O. Gunnarsson, PRB 70, 224518.

## Doping dependence of softening



## Line width in HF solution

HF softening OK for  $\delta \sim 0.10 - 0.15$ , but width much smaller than exp.  
Real and imaginary part of self-energy related. Contradiction?  
 $t$ - $J$  calculation for  $4 \times 4$  cluster (details of  $Im \Pi$  not reliable).  
• Linewidth (0.4 meV) smaller than in  $t$ - $J$  model (3 meV) (many low-lying excitations in  $t$ - $J$  model).  
• Full-breathing mode softens more than half-breathing mode, contrary to exp.  
O. Rösch and O. Gunnarsson, PRB 70, 224518.

## Summary

- LDA suggests very weak coupling to half-breathing phonon.
- Many-body treatment ( $t$ - $J$  model) suggests much stronger coupling.
- Experimental evidence for substantial coupling.
- LDA fails to describe dispersion of half-breathing phonon.
- $t$ - $J$  model gives proper dispersion and much larger width.